

# *Conformal Fixed Point of $SU(3)$ Gauge Theory with 12 Fundamental Fermions in the Twisted Polyakov Loop Scheme*

Eigo Shintani (RIKEN-BNL)

and

T. Aoyama(KMI), H. Ikeda(KEK), E. Itou(Osaka), M.  
Kurachi(KMI), C.-J. D. Lin(NCTU&NCTS), H.  
Matsufuru(KEK), K. Ogawa(NCTU), H. Ohki(KMI), T.  
Onogi(Osaka) and T. Yamazaki(KMI)

arXiv:1109.5806[hep-lat]

# Plan of my talk

- Introduction
- Twisted Polyakov loop (TPL) scheme
- $N_f=12$  results
- Summary and outlook

# Introduction

# Motivation

- **Search of non-trivial IR fixed point (IRFP) in  $N_f$  flavor QCD**

- For dynamical EW breaking in (walking) technicolor model

- One of the candidates of BSM (revival from 80's...)

Holdom (1986), Yamawaki (1986)

Appelquist (1986), ...

- Predict dynamical symmetry breaking in **strongly interacting**

**gauge theory**, analogy with QCD

- To get rid of large FCNC, S-parameter and small quark mass

1. Nearly conformal theory

2. Large anomalous mass dimension

Does the naïve extension of QCD to large flavor theory realize the above conditions ?  $\rightarrow$  **non-perturbative study is important !**

- For phase structure

Chiral broken phase and confinement phase will be different or not ?

$\rightarrow$  perception of phase structure of QCD

Gardi and Karliner (1998),

Miransky and Yamawaki (1997),

Ryttov and Sannino (2008)



©Tezuka Productions

# Motivation

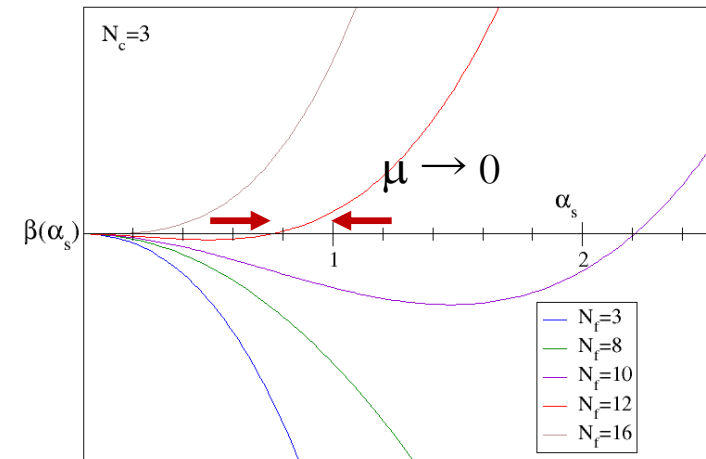
## • Caswell-Banks-Zaks IRFP

In 2-loop perturbation, beta function

$$\beta(g^2) \equiv \frac{\partial \alpha}{\partial \ln \mu^2} = -\beta_0 \alpha_s^2 - \beta_1 \alpha_s^3, \quad \alpha_s = g^2/(4\pi)$$

$$\beta_0 = \frac{1}{4\pi} \left( \frac{11}{3} N_c - \frac{2}{3} N_f \right),$$

$$\beta_1 = \frac{1}{16\pi^2} \left( \frac{34}{3} N_c - \left\{ \frac{13}{3} N_c - \frac{1}{N_c} \right\} N_f \right)$$



IRFP ( $N_c=3$ ) in  $N_f^{\text{cr}} (\simeq 8) < N_f < N_f^{\text{af}} (\simeq 17)$  Caswell (1974), Banks and Zaks (1982)

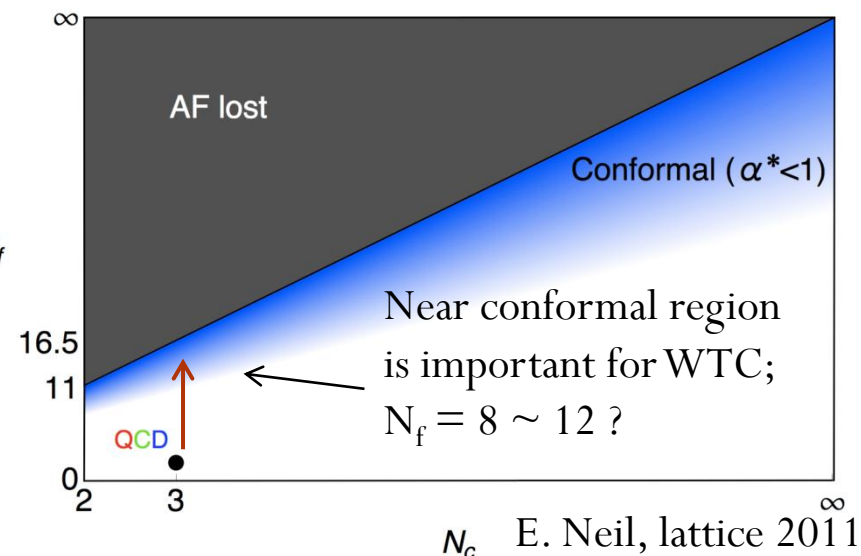
- $m = 0$  (fundamental rep.)
- $g(\mu \rightarrow 0) = g^*$ , IR conformal

$\Rightarrow$  “conformal window”

Chiral symmetry will not be broken,  $N_f$

and 1<sup>st</sup> order transition occurs ?

This perception will be clear in the lattice calculation.



# Motivation

## • Lattice study in 12 flavor

### • Appelquist, et al. (2008,2009)

- Measurement of the running coupling in **SF scheme** in staggered fermion
- See plateau region in  $g^2 \simeq 4-5$  that is IRFP

$\Rightarrow \underline{N_f = 12 \text{ is conformal}}$

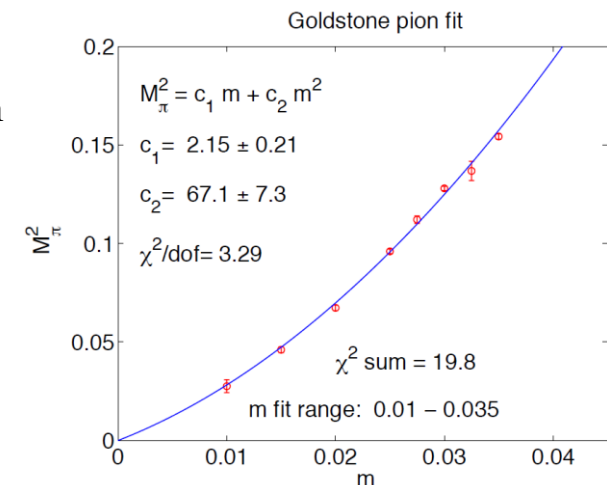
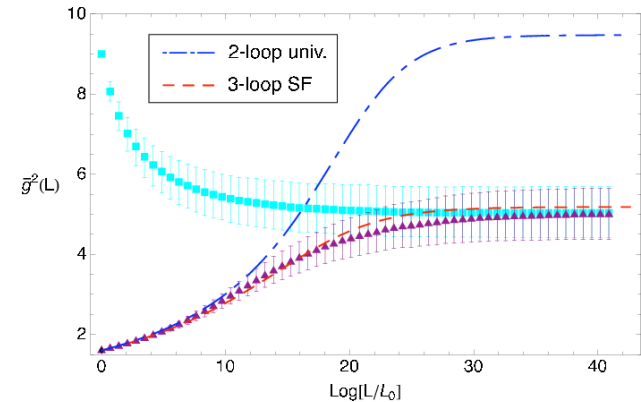
### • Fodor et al. (2009,2011)

- Goldstone pion and composite hadron spectrum in the improved staggered fermion
- Goldstone pion behaves like  $\chi$ SB, and  $\langle \bar{\psi}\psi \rangle \neq 0$

$\Rightarrow \underline{N_f = 12 \text{ is in broken phase}}$

### • A. Hasenfratz, (2010,2011)

- Monte Carlo renormalization group (MCRG)
- Beta function will cross zero from negative to positive  $\Rightarrow \underline{N_f = 12 \text{ is conformal}}$



Need confirmation in 12 flavor with different scheme

# New scheme

- **Twisted Polyakov loop scheme**

- Different lattice scheme from SF and MCRG
- Avoid a fake (unphysical) fixed point
- Free from  $O(a)$  discretization error
- Practically cheaper cost than Wilson loop scheme      Bilgici et al. (2009),  
Not need large volume, zero fermion mass simulation      Holland (2009)
- Consistency check with Wilson loop scheme in quenched QCD has been already done.      Itou (2010)
- Measure the renormalized running coupling **in the continuum limit**  
Obtain anomalous dimension of coupling constant (universal object)
- First target on 12 flavor QCD in fundamental representation

# Twisted Polyakov loop scheme



# Polyakov loop in twisted boundary

- **Twisted boundary condition**

- To avoid degeneracy of vacua (“toron”) ‘tHooft (1979), Gonzalez-Arroyo (1988)

- Suppression of  $Z(3)$  phase transition

⇒ well defined perturbative expansion of Wilson loop Divitiis, et al. (1994)

For gauge field (link variable):

$$U_\mu(x + \hat{\nu}L/a) = \Omega_\nu U_\mu(x) \Omega_\nu^\dagger, \quad \Omega_\nu: \text{twist matrix}$$

$$\Omega_1 \Omega_2 = e^{i2\pi/3} \Omega_2 \Omega_1, \quad \Omega_\nu \Omega_\nu^\dagger = 1, \quad \Omega_\nu^3 = 1, \quad \text{Tr } \Omega_\nu = 0$$

For fermion field:

$$\psi \begin{matrix} \text{smell} & \text{color} \\ \underbrace{a}_{\alpha} & \underbrace{b} \end{matrix} (x + \hat{\nu}a/L) = e^{i\pi/3} (\Omega_\nu)^{aa'} \psi_{\alpha}^{a'b'}(x) (\Omega_\nu^\dagger)^{b'b}$$

spinor

Parisi, (1983)

Wong, et al. (2006)

Add “smell” degree of freedom being same as color, corresponding to extra flavor, in order to avoid inconsistency with translational invariance.

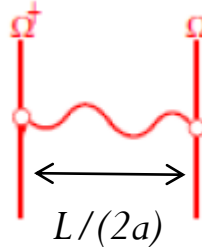
Using staggered fermion, flavor number is  $N_f = 4 \times 3 = 12$

# Polyakov loop in twisted boundary

- Ratio of Polyakov loop correlator

$$P_x(y, z, t) = \text{Tr} \left[ \prod_x U_x(x, y, z, t) \underbrace{\Omega_x}_{\text{Gauge inv.}} \underbrace{e^{i2\pi/3}}_{\text{Trans. inv.}} \right], \quad P_z(x, y, t) = \text{Tr} \left[ \prod_z U_z(x, y, z, t) \right]$$

$$\langle P_x(t = L/(2a)) P_x^*(0) \rangle$$

$$= \text{Diagram} \sim k g_0^2 + \mathcal{O}(g_0^4)$$


$$\langle P_z(t = L/(2a)) P_z^*(0) \rangle$$

$$= \text{Diagram} \sim 1 + \mathcal{O}(g_0^4)$$


Divitiis (1994)

$$(g_{\text{TP}}^{\text{lat}})^2 = \frac{1}{k} \frac{\langle \sum_{y,z} P_x(y, z, L/(2a)) P_x^*(0, 0, 0) \rangle}{\langle \sum_{x,y} P_z(x, y, L/(2a)) P_z^*(0, 0, 0) \rangle}, \quad (g_{\text{TP}}^{\text{lat}})^2|_{\text{tree}} = g_0^2$$

$$k = 0.0318471147 \dots + 0.00453(a/L)^2$$

- $(g_{\text{TP}}^{\text{lat}})^2$  starts from log divergence (renormalized), no linear divergence due to cancelation between numerator and denominator.

→ free from  $O(a/L)$  discretization error

# Polyakov loop in twisted boundary

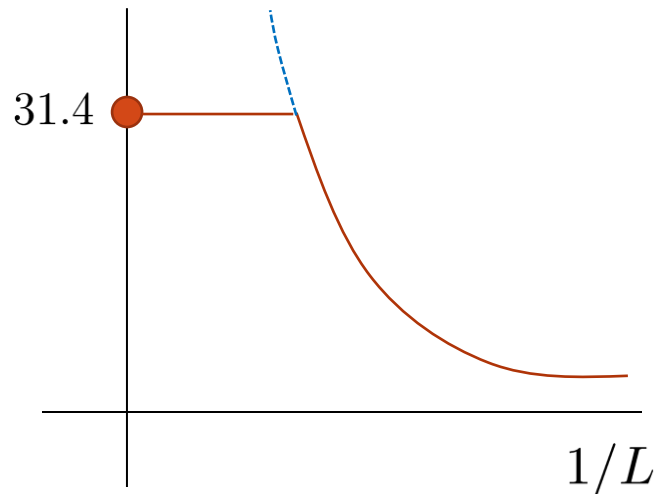
- **Behavior in IR region**

- In the infinite size, Polyakov loop correlator does not depend on boundary, thanks to cluster theorem, then

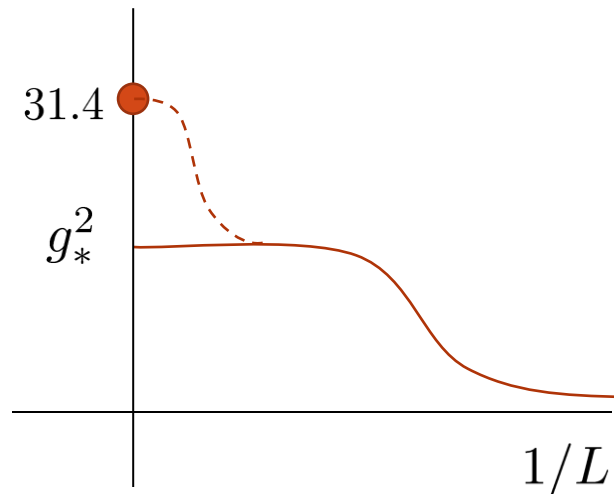
$$\lim_{L \rightarrow \infty} \frac{\langle \sum_{y,z} P_x(y, z, L/(2a)) P_x^*(0, 0, 0) \rangle}{\langle \sum_{x,y} P_z(x, y, L/(2a)) P_z^*(0, 0, 0) \rangle} = 1 \quad \Rightarrow \quad [g_{\text{TPL}}^2]_{\text{max}} \sim \frac{1}{k} \simeq 31.4$$

- The above value is maximum in the IR limit. It turns out that it is able to distinguish IRFP from other (e.g. QCD-like) theory.

QCD-like (no-IRFP or  $g_*^2 \geq 31.4$ )



IRFP (or walking)



# Running coupling constant

- **Step scaling method**

- Inverse of volume is interpreted as a scale.
- Changing the volume gives the stepping behavior of running coupling
- Start from reference point:  $u_0$   
which is weak coupling constant in the continuum theory
- Set a step parameter:  $s$   
Large step parameter is recommended to investigate deep IR region and suppress correlation

- **Step scaling process:**

1.  $u_0 = g_{\text{TPL}}^2(1/L),$
2.  $(g_{\text{TPL}}^{\text{lat}})^2(a, a/L) = u_0, \quad \Sigma(a, a/(sL)) = (g_{\text{TPL}}^{\text{lat}})^2(a, a/(sL))$
3.  $\sigma(u_0) = \lim_{a \rightarrow 0} \Sigma(a/(sL), a),$
4.  $u_1 = \sigma(u_0)$  return to 1.

e.g.  $s = 2, aL = 4, 6, 8 \rightarrow saL = 8, 12, 16$ , and taking continuum limit with several beta simulation, we can get step scaling **in the continuum theory**.

- If an IRFP exists in  $u=u^*$ , see  $\sigma(u^*)/u^* = 1$

Nf=12 results

# Nf=12, staggered fermion

- Dynamical fermion with “smell” in HMC  
4 flavor(default)  $\times$  3 smell = 12 flavor effectively
- Simulation on the **massless** point, thanks to twisted boundary  
No need of chiral extrapolation

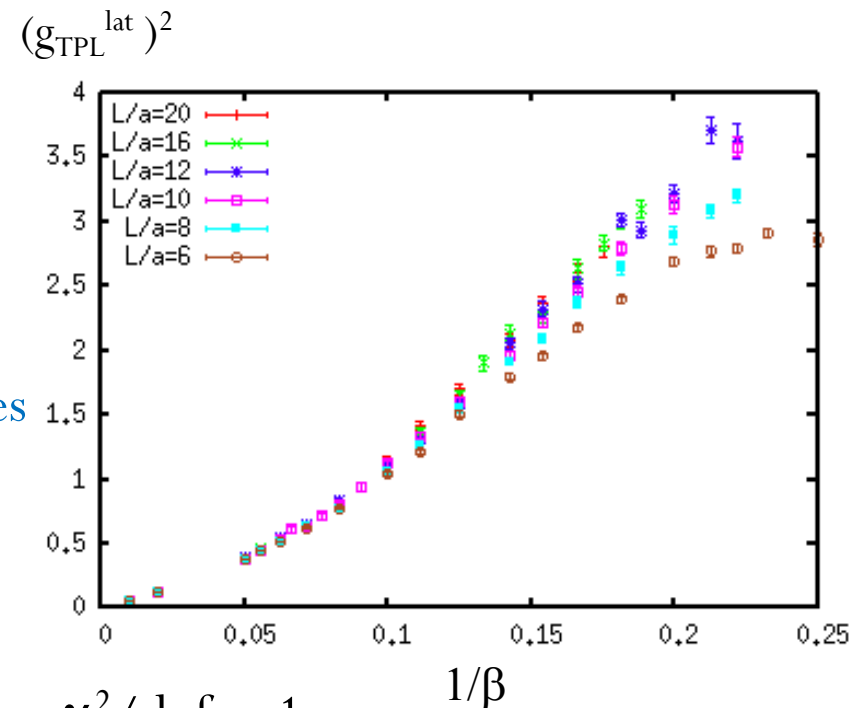
- Parameters:

- 4 lattice size,  $L/a = 6, 8, 10, 12$   
and  $sL/a = 9, 12, 15, 18$  with  $s = 1.5$   
Interpolating points
- $1/\beta = 0.01 \text{ -- } 0.25$
- 8k(high  $\beta$ ) – 897k(low  $\beta$ ) trajectories

- Fit function for interpolation

$$(g_{\text{TPL}}^{\text{lat}})^2 = \frac{6}{\beta} + \sum_{i=1}^N C_i(a/L)\beta^{-i-1}$$

with  $N = 3 - 5$  depending on lattice size,  $\chi^2/\text{dof} \sim 1$



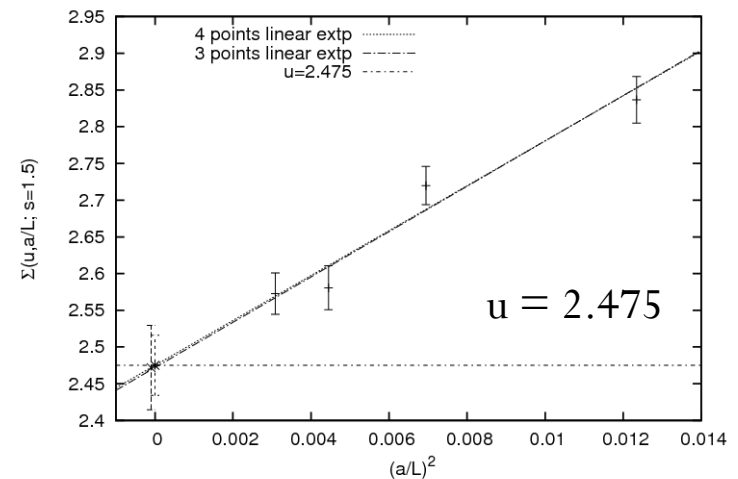
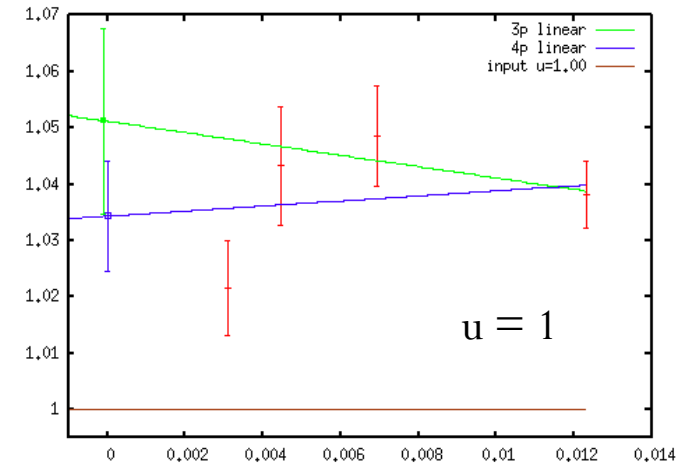
# Nf=12, staggered fermion

- **Continuum limit**

- After s step, we need to take the continuum limit.

In TPL scheme,  $O(a/L)$  error is absent.

- We perform a linear extrapolation for  $(a/L)^2$  using data in  $sL/a = 9, 12, 15, 18$ .
  - Linear function well describes lattice data.
  - 3 and 4 points fit are not so different. It turns out that systematic error of scaling violation is not so large.



# Nf=12, staggered fermion

## Running coupling constant in TPL scheme

- Behavior of ratio,  $\sigma(u)/u$  vs  $u$ 
  - From  $u \simeq 0.5$ ,  $\sigma(u)/u$  grows up rather than 2-loop, and around  $u \simeq 1.3$  it declines to 1.
  - $u \simeq 2.5$  this ratio crosses 1, and it will be conformal point.
  - Anomalous dimension of coupling  $\gamma$ , which is slope at  $u^*$

$$\gamma = 0.52^{+27}_{-24} \quad \begin{array}{l} \text{2-loop: } \gamma \sim 0.36 \\ \text{SF scheme: } \gamma = 0.13(3) \end{array}$$

consistent with 2-loop result, but error is still large.

$\Rightarrow$  future works

## Running behavior

UV region: Consistent with 2-loop

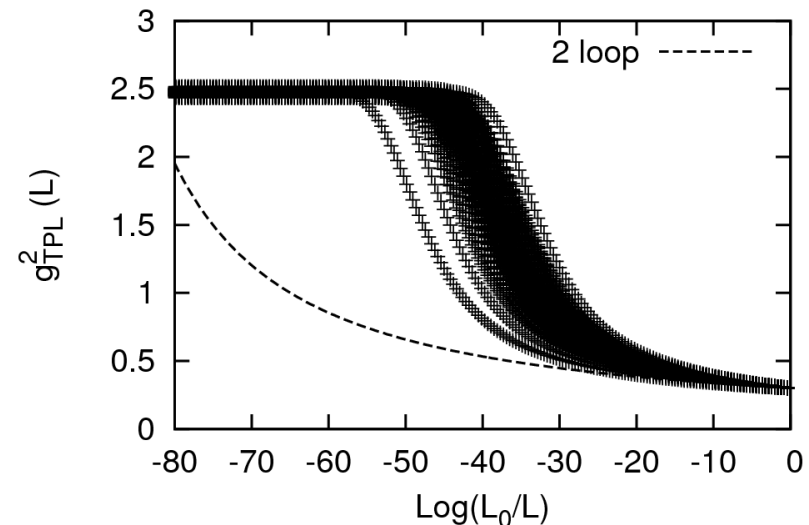
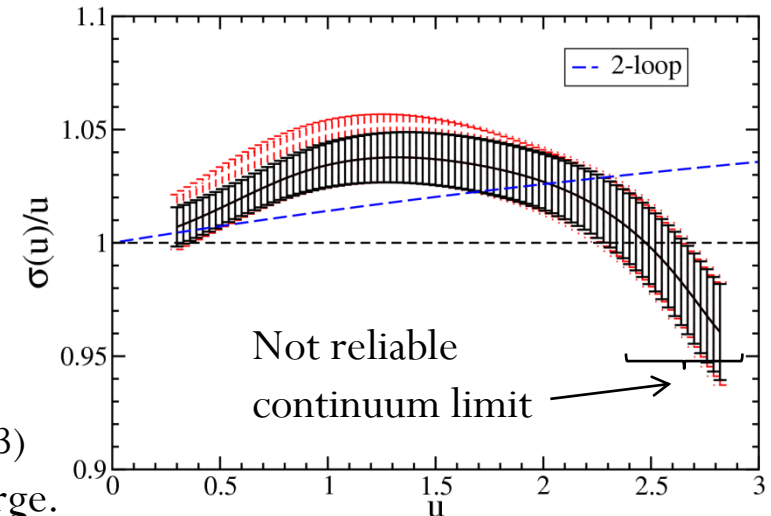
IR region:

Observe a unique plateau region for every distribution of JK ensembles, and then

$$(g_{\text{TPL}}^2)^* = 2.48(18)(^{+7}_{-8}) \ll [g_{\text{TPL}}^2]_{\text{max}}$$

Black solid : statistical error

Red dashed: total error





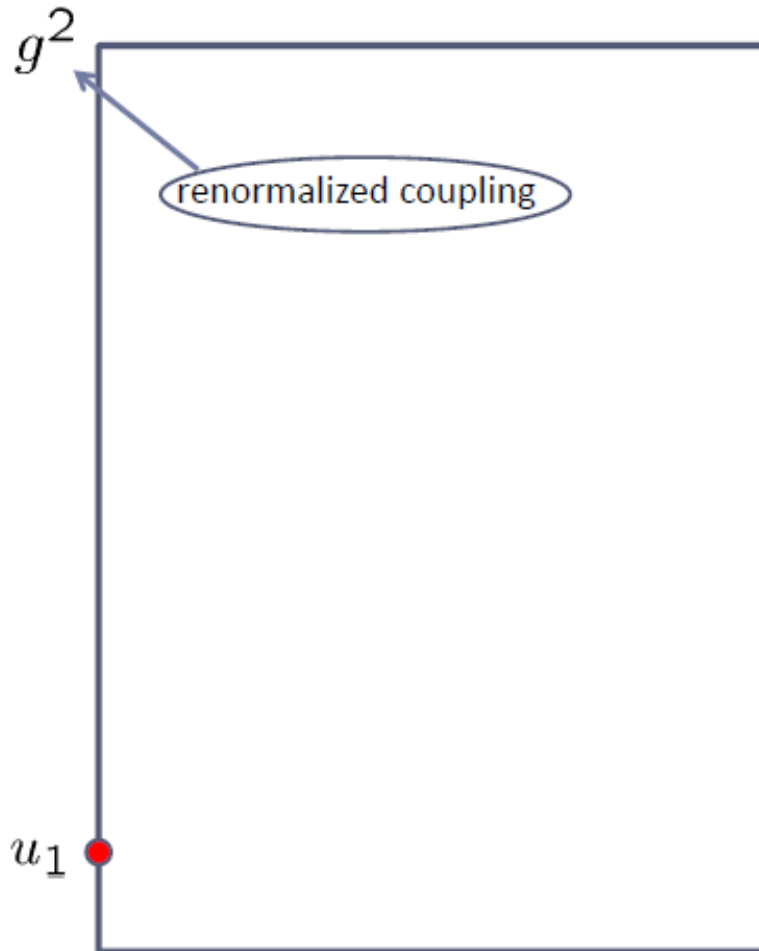
# Summary and outlook

- Simulation of  $N_f = 12$  SU(3) gauge with fundamental fermion using “Twisted Polyakov loop” scheme.
- High statistics and rigorous lattice calculation at **massless point**
  - Take **the continuum limit**, and control scaling violation
- Running coupling constant slows down and stops in  $(g_{\text{TPL}}^2)^* = 2.48(18)(^{+7}_{-8})$
- **Our results establish that**  
 **$N_f = 12$  QCD (fund.) has an infrared conformality.**
  - Appelquist, et al.
  - Haesenfratz
  - Aoki (KMI)
  - Pallente, et al.
- On-going works:
  - Mass anomalous dimension in this scheme  
Obtained from correlator of Goldstone pion
  - SU(2) fundamental in  $N_f = 8$ , this is almost done

# Backup

# Example

## Step scaling



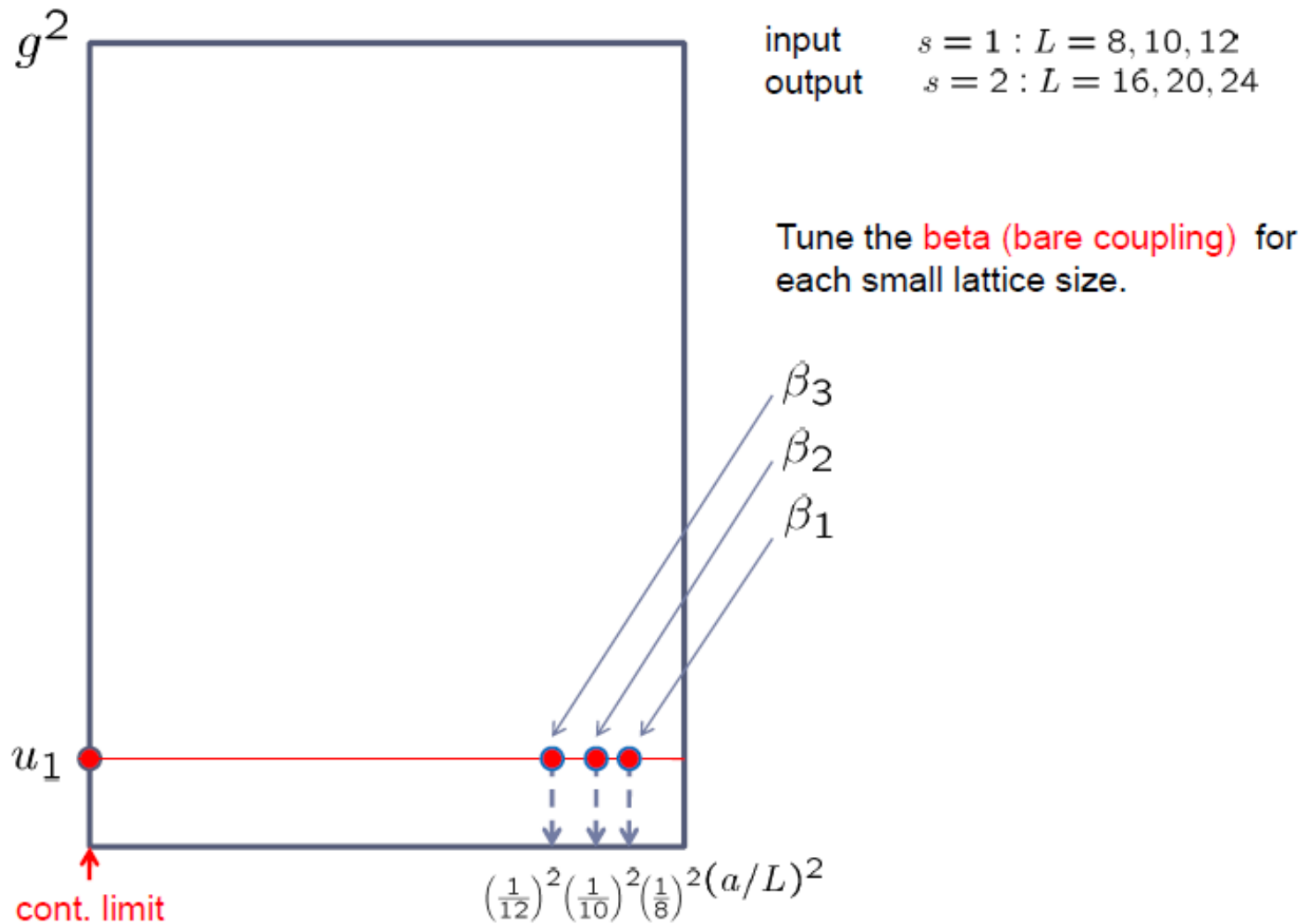
input  $s = 1 : L = 8, 10, 12$   
output  $s = 2 : L = 16, 20, 24$

Choose a value of the  
renormalized coupling constant  
at energy scale ( $\mu = 1/L_0$ )

$(a/L)^2 \leftarrow$  inverse of lattice size

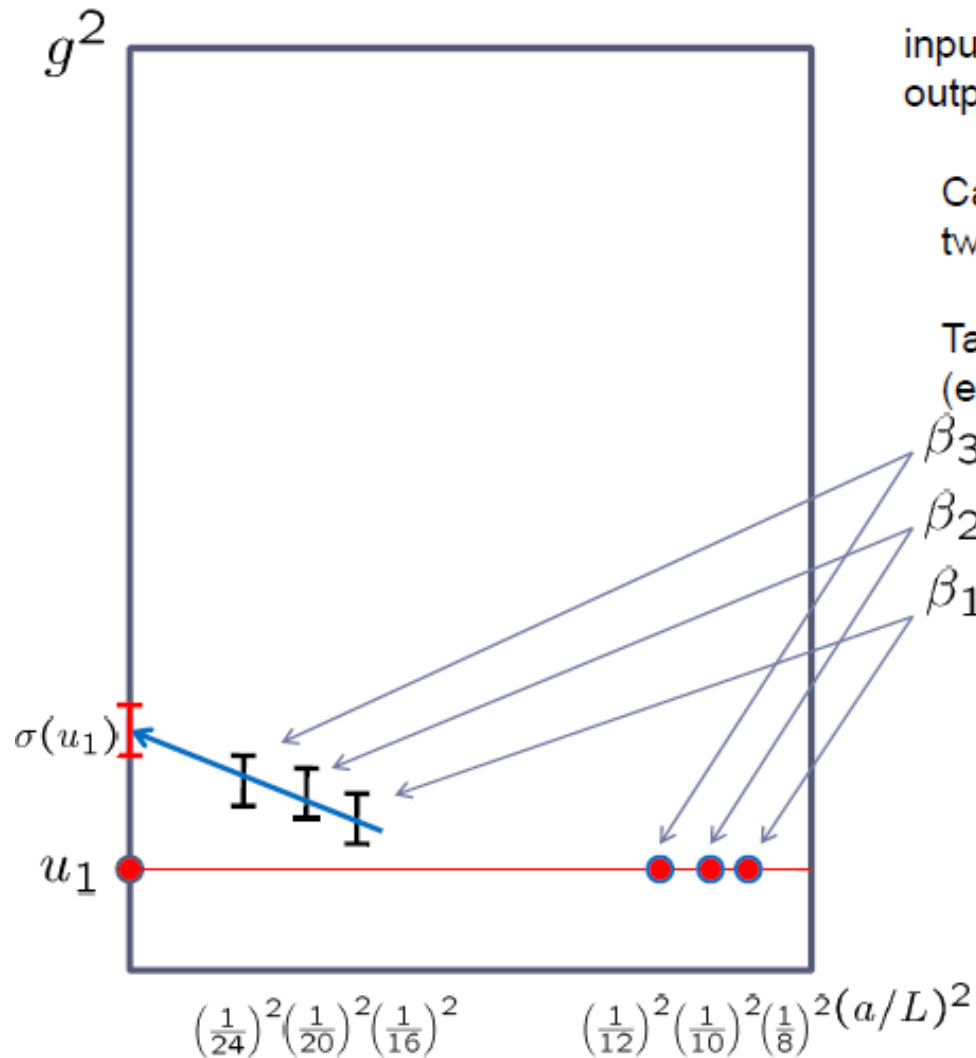
# Example

## Step scaling



# Example

## Step scaling



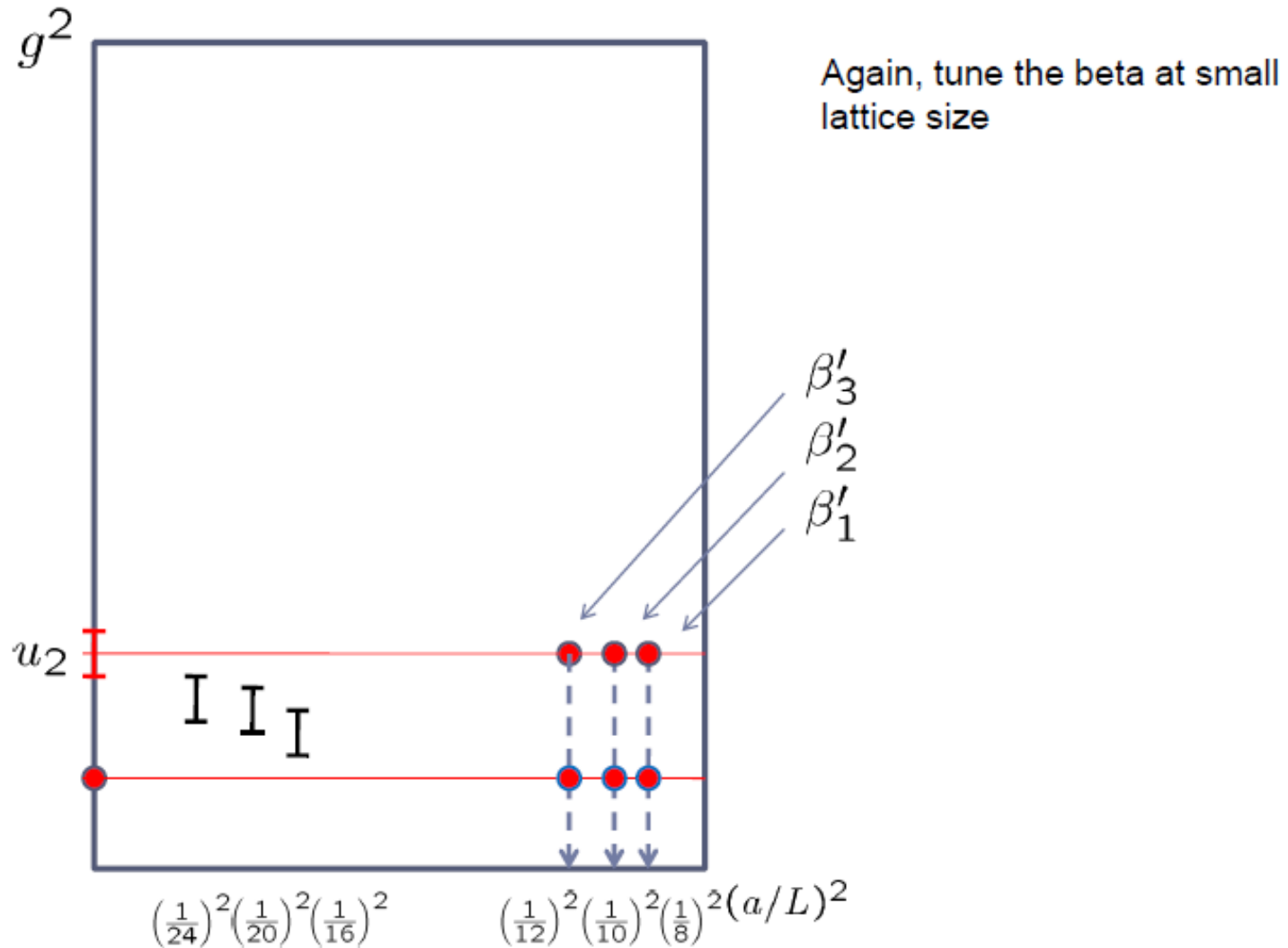
input  $s = 1 : L = 8, 10, 12$   
output  $s = 2 : L = 16, 20, 24$

Carry out the simulation for the twice size of lattice.

Take the continuum limit  
(energy scale  $\mu = 1/2L_0$ )

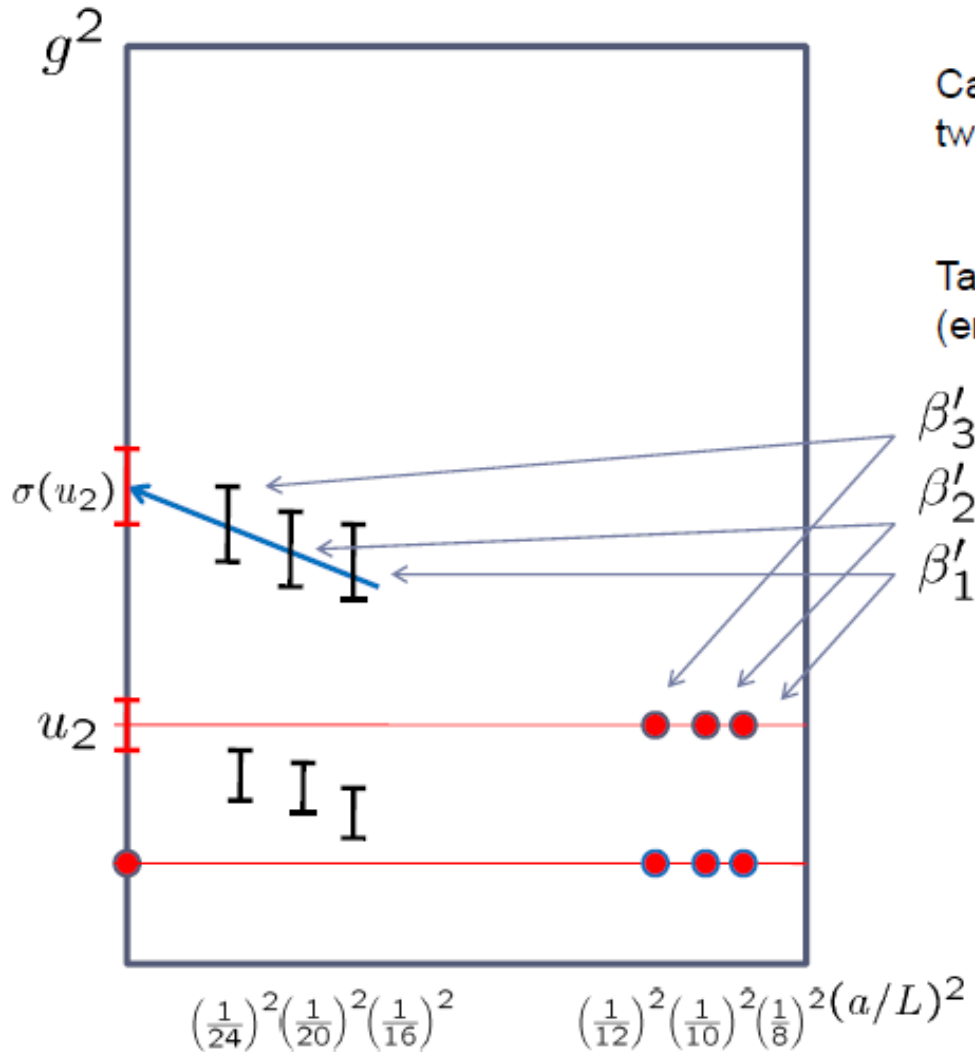
# Example

## Step scaling



# Example

## Step scaling

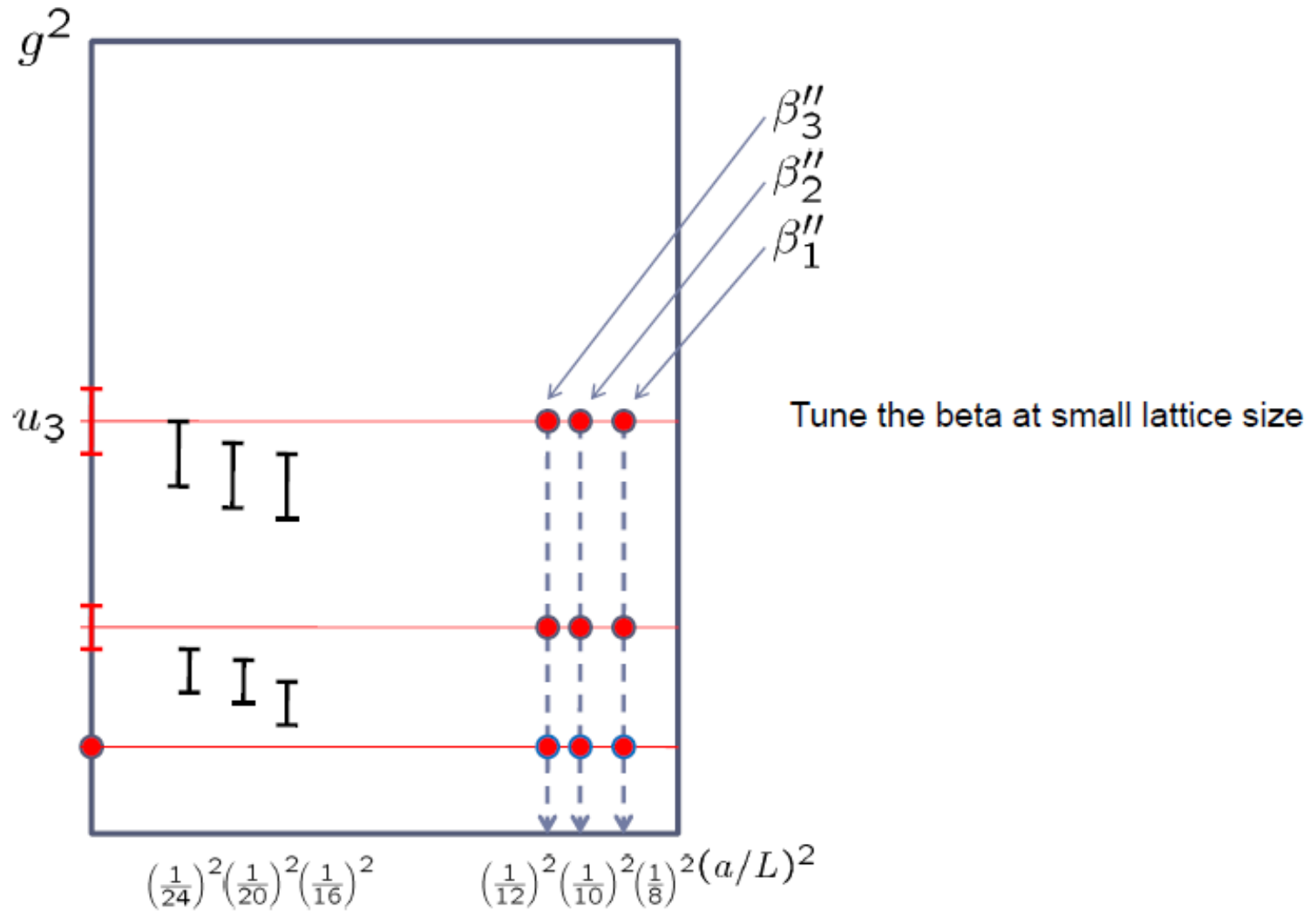


Carry out the simulation for the twice size of lattice.

Take the continuum limit  
(energy scale  $\mu = 1/4L_0$ )

# Example

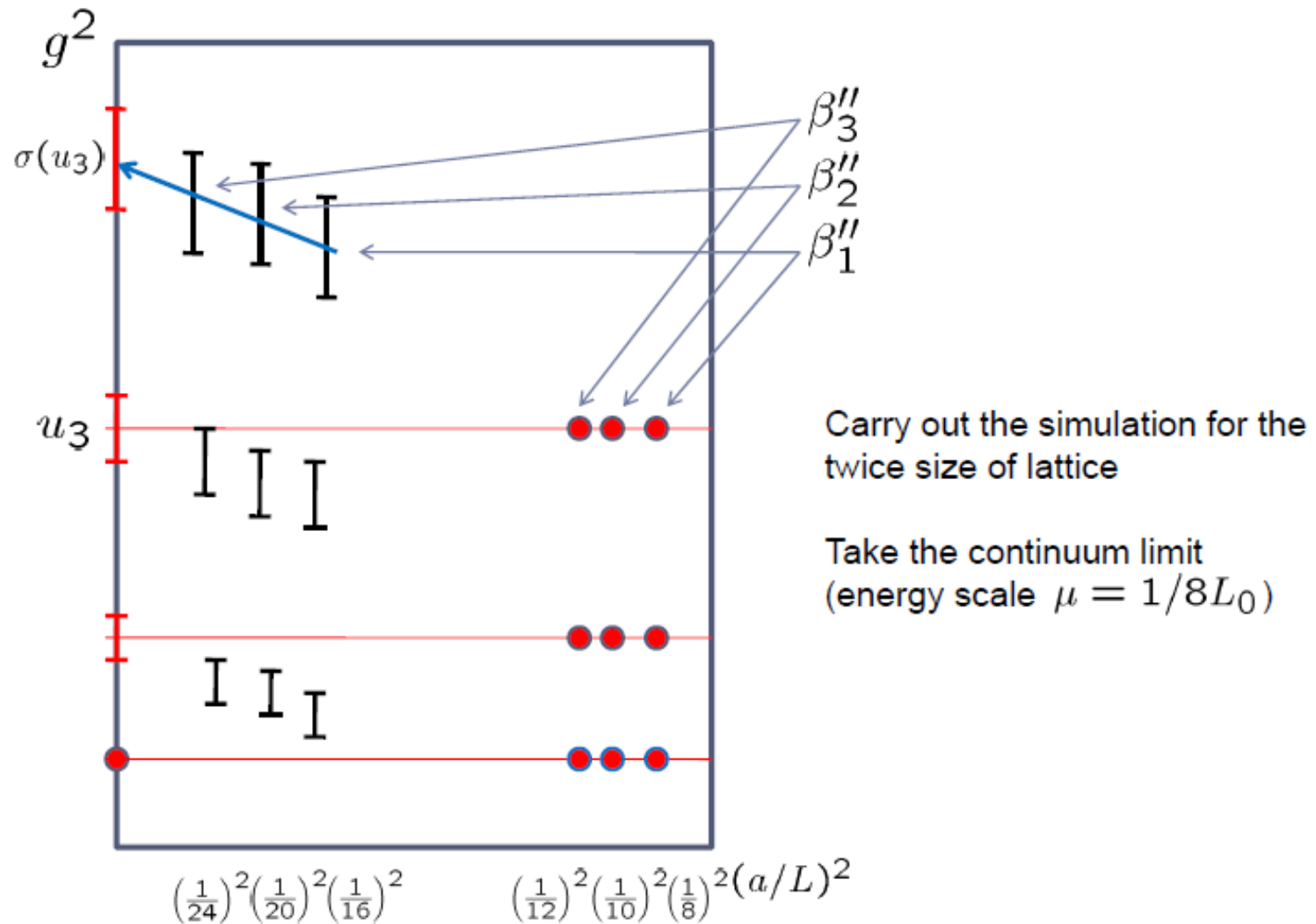
## Step scaling





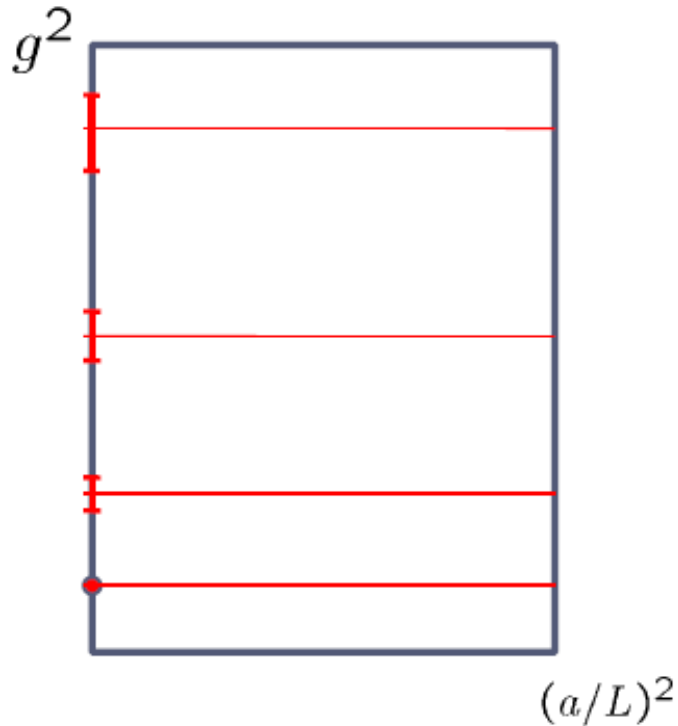
# Example

## Step scaling



# Example

## Step scaling



We obtain the scaling of the running coupling.

